

# Nuclear diagnostics and Magnetic Resonance Imaging

## Lecture 10: Magnetic Resonance Imaging: spatial localisation

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# Outline

## 1 Spatial localisation

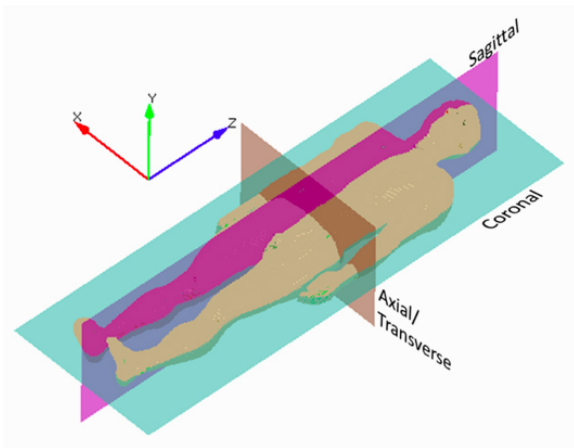
- Introduction
- Slice selective excitation
- Encoding spacial information
- Encoding spatial information into net magnetisation

## 2 Lecture summary

## Section 1

# Spatial localisation

# Introduction



Conventional terminology & orientation of RH coordinate system

Contrast between tissues is afforded by RF  $B_1$  pulse sequences such as those discussed in the preceding lectures

To make an image, need to localise the signals to appropriately small regions of space

To localise signals exploit:

- Resonance, i.e. Larmor frequency  $\nu = \gamma B$
- By making  $B$  a function of position

i.e. make  $\nu$  a function of position:

$$\nu(x, y, z) = \gamma B(x, y, z)$$

# Slice selective excitation

Goal: excite a slice of tissue of thickness  $\delta$

So far a uniform “main field”  $\mathbf{B}_0 = B_0 \hat{\mathbf{k}}$  has been considered

Require to make  $B_z$  a function of position to make Larmor frequency position dependent

Apply “gradient” fields  $G_i$  such that  $B_z$  becomes:

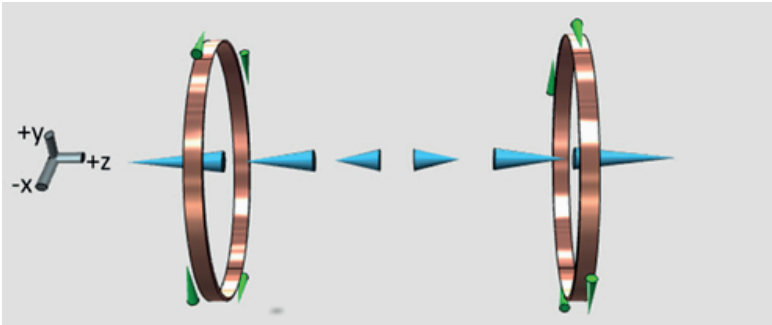
$$B_z(x, y, z, t) = B_0 + xG_x(t) + yG_y(t) + zG_z(t)$$

Ideally  $G_i$  only have one field component directed along the  $z$  direction so that:

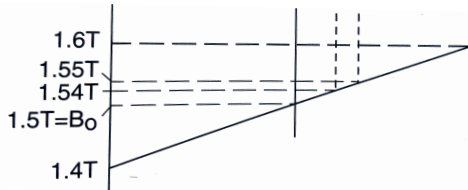
$$\mathbf{B} = B_z(x, y, z, t) \hat{\mathbf{k}}$$

With appropriate choice of  $G_i$  can generate a field gradient in any direction

## Transverse slice; i.e. plane at fixed $z$



Example:  
Helmholtz coils in  
opposition



Ideal gradient:  
 $G_z = \text{constant}$

## Transverse slice; slice thickness and bandwidth

Lets say that response needs to be isolated to a slice:  $\delta z = 5 \text{ mm}$  centred about  $z = 0$

Take:

- The magnitude of the main field to be  $B_0 = 1.5 \text{ T}$
- The field gradient  $G_z = 50 \text{ mT m}^{-1}$
- $\gamma = 42.58 \text{ MHz T}^{-1}$

Take the slice to be  $-2.5 < z < 2.5 \text{ mm}$ , then the Larmor frequency will run over the following range:

$$\nu_{\min} = (1.5 - 0.125 \times 10^{-3}) \times 42.58 \approx 63.8646 \text{ MHz}$$

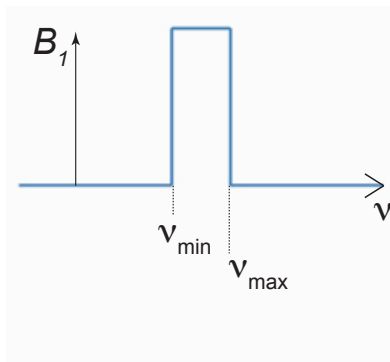
$$\nu = 1.5 \times 42.58 \approx 63.87 \text{ MHz}$$

$$\nu_{\max} = (1.5 + 0.125 \times 10^{-3}) \times 42.58 \approx 63.8753 \text{ MHz}$$

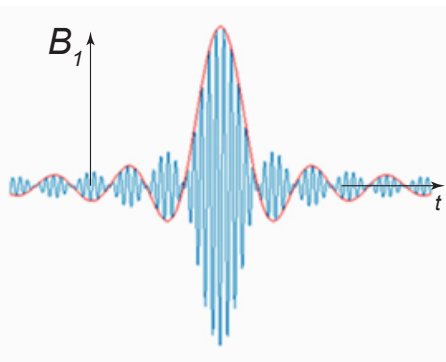
So, the spread of frequencies, the **bandwidth**,  $\Delta\nu$  is:

$$\Delta\nu = 63.8646 - 63.8753 \approx 10.7 \text{ kHz}$$

# Transverse slice; excitation of spins in slice



Idealised, square frequency distribution

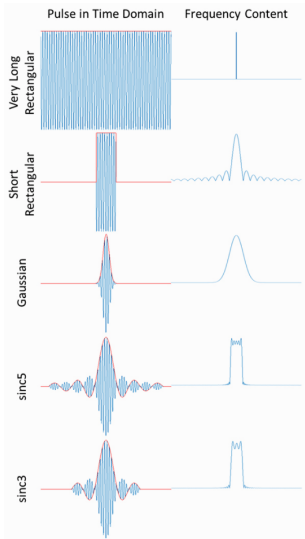


Fourier transform of square frequency distribution

$B_1$  oscillates at  $\nu$ , amplitude is modulated according to “sinc” function (red line)



# Transverse slice: excitation pulses



Frequency content of a variety of excitation pulses:

- *Very long rectangular*: narrow band of Larmor frequencies
- *Short rectangular*: frequency distribution follows “sinc” function:

$$A(\nu) \propto \text{sinc}(\nu) = \frac{\sin \nu}{\nu}$$

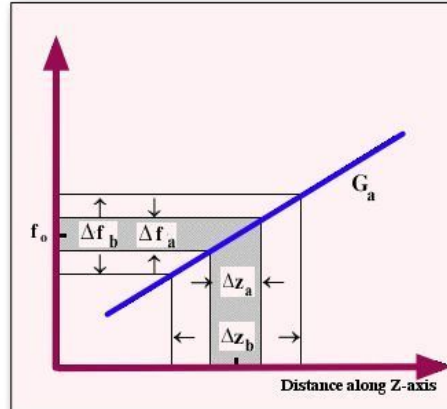
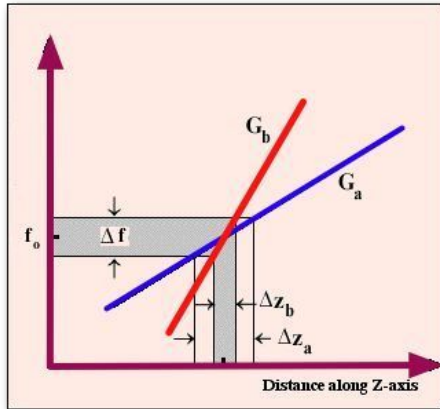
where  $A(\nu)$  is the amplitude of contribution at frequency  $\nu$

- *Gaussian*: Fourier transform of Gaussian in  $t$  is a Gaussian in  $\nu$
- *sincN*: Since square pulse requires contributions over all  $\nu$ , the frequency range is often truncated. The “sincN” function represents a sinc function for which the frequency range is truncated after N zero crossings

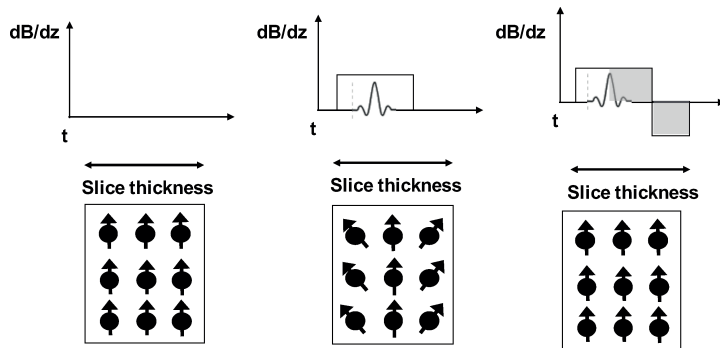
# Transverse slice: determining the slice thickness

Slice thickness is determined by bandwidth ( $\Delta\nu$ ) and field gradient ( $G_z$ )

Sorry for the change in notation!



# Transverse slice: spin rephasing pulse



Larmor frequency across slice changes. So, over the time that the gradient pulse is applied, the spins precess at different rates

Therefore, at the end of the pulse the phase of the spins differs as a function of  $z$

A rephasing pulse which reverses the field gradient (i.e. for which  $G_z \rightarrow -G_z$ ) is applied

## Transverse slice: spin rephasing pulse

Size of the spin rephasing pulse is determined by considering the rate at which the phase difference accumulates

Rate of precession is given by the Larmor frequency,  $\omega$ , so change in phase of a spin during the gradient pulse is given by:

$$\Phi = \omega\tau = \gamma(B_0 + zG_z)\tau$$

where  $\tau$  is the length of the gradient pulse in time

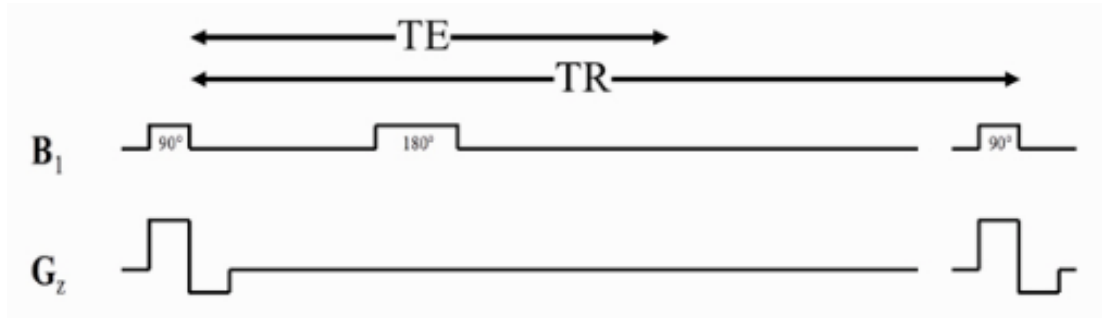
So, phase difference between edges of the slice and the centre is:

$$\Delta\Phi = \gamma\tau G_z \frac{\delta z}{2}$$

So, rephasing pulse,  $G_z^{\text{rephase}}$ , and the length over which it is applied,  $\tau^{\text{rephase}}$  must satisfy:

$$G_z^{\text{rephase}} \tau^{\text{rephase}} = G_z \frac{\tau}{2}$$

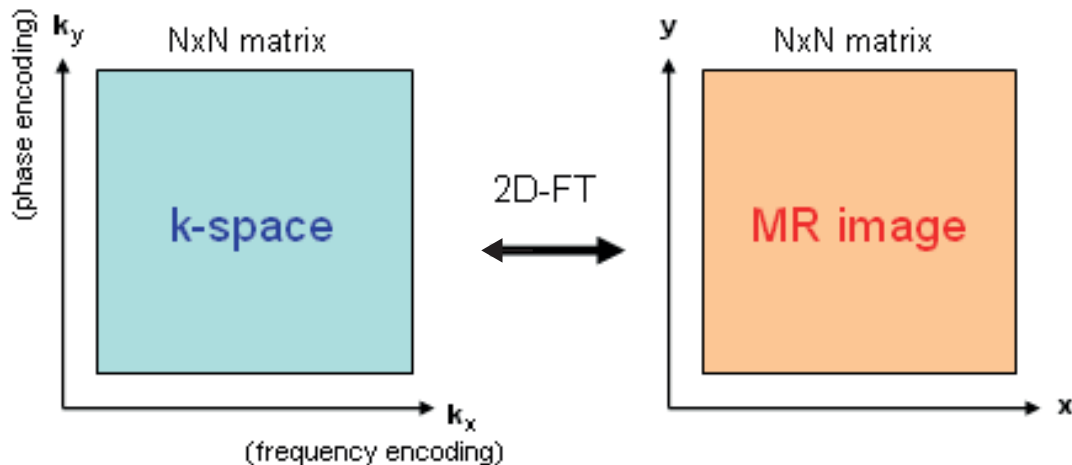
## Transverse slice: partial spin-echo pulse sequence



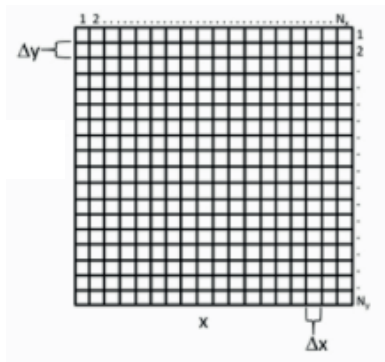
$B_1$  rotates net magnetisation in the selected slice with gradient pulse applied

# Encoding spatial information into the net magnetisation

The basis is a 2D Fourier transform:



## 2D Fourier transformation



2D image in “coordinate space”,  $x, y$ , presented in pixel grid

Field of view, FOV, in coordinate space:

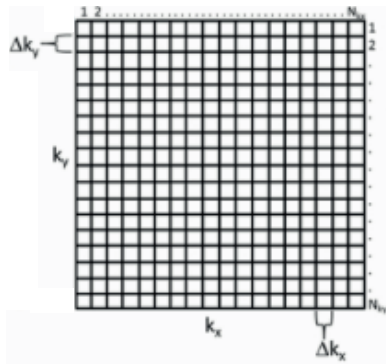
$$(x_{\max} - x_{\min}, y_{\max} - y_{\min})$$

Pixel size (resolution):

$$\Delta x = \frac{x_{\max} - x_{\min}}{N_x}$$

$$\Delta y = \frac{y_{\max} - y_{\min}}{N_y}$$

## 2D Fourier transformation



2D image in “ $k$  space”,  $k_x, k_y$ , presented in pixel grid

Field of view, FOV, in  $k$  space:

$$(k_{x \max} - k_{x \min}, k_{y \max} - k_{y \min})$$

Pixel size (resolution):

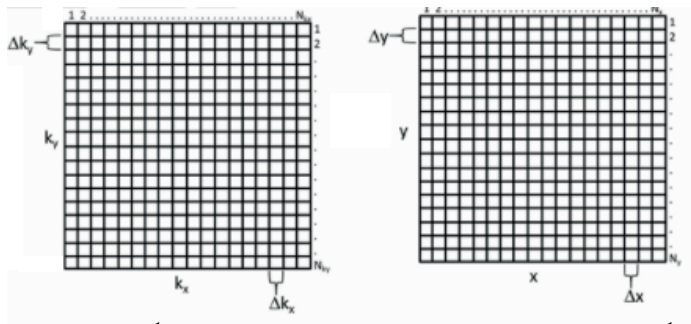
$$\Delta k_x = \frac{k_{x \max} - k_{x \min}}{N_x}$$

$$\Delta k_y = \frac{k_{y \max} - k_{y \min}}{N_y}$$



## 2D Fourier transformation

Transformation between resolution in coordinate-space and  $k$ -space representations:



$$\Delta k_x = \frac{1}{(x_{\max} - x_{\min})}$$

$$\Delta k_y = \frac{1}{(y_{\max} - y_{\min})}$$

$$\Delta x = \frac{1}{(k_{x \max} - k_{x \min})}$$

$$\Delta y = \frac{1}{(k_{y \max} - k_{y \min})}$$

## 2D Fourier transformation

Define  $\rho(x, y)$  to be the intensity pixel-by-pixel in coordinate space.

2D Fourier transform from coordinate to  $k$  space is then:

$$S(k_x, k_y) = \int_{y_{\min}}^{y_{\max}} \int_{x_{\min}}^{x_{\max}} \rho(x, y) \exp(-i2\pi k_x x) \exp(-i2\pi k_y y) dx dy$$

where  $S(k_x, k_y)$  is the intensity pixel-by-pixel in  $k$  space

Inverse Fourier transform takes  $k$ -space intensity map to coordinate-space intensity map:

$$\rho(x, y) = \int_{k_y \min}^{k_y \max} \int_{k_x \min}^{k_x \max} S(k_x, k_y) \exp(i2\pi k_x x) \exp(i2\pi k_y y) dk_x dk_y$$

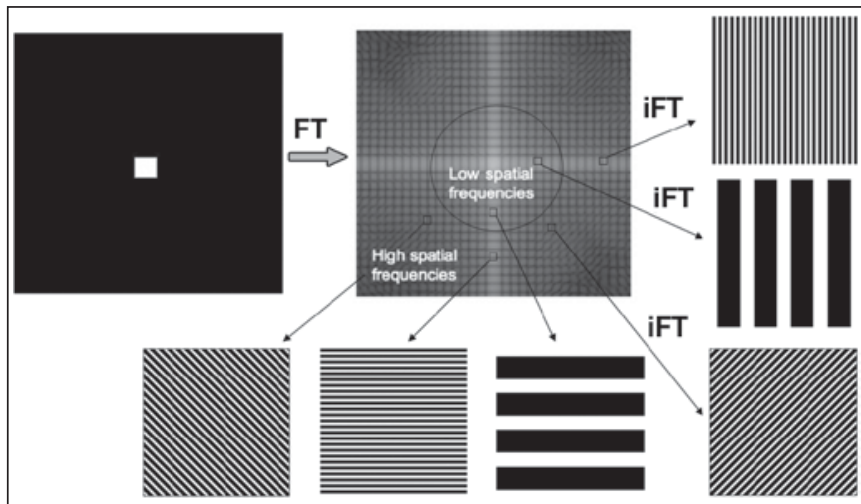
## Example one: a single dot



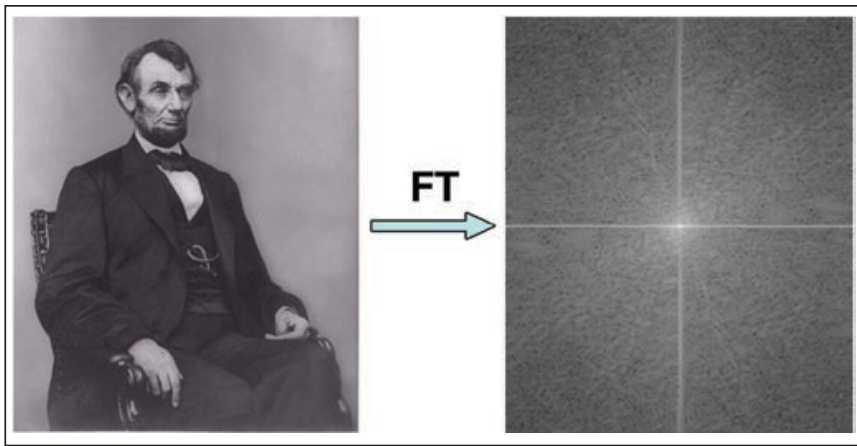
## Example two: three dots



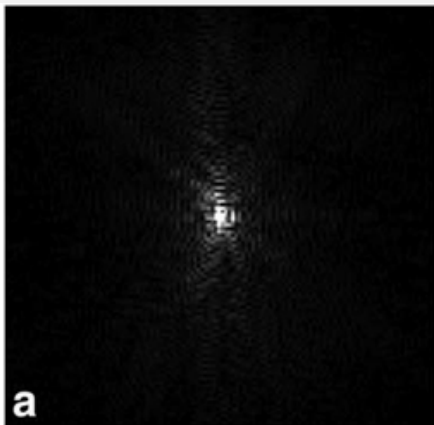
## Example three: Square in centre of field of view



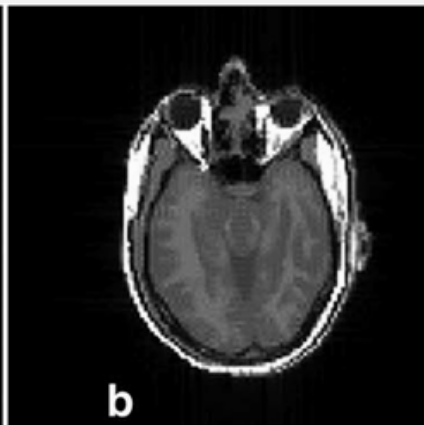
## Example three: Abraham Lincoln



## Example three: Abraham Lincoln



(a)  $k$ -space image of head



(b) coordinate-space image of head

Challenge: record  $k$ -space image using NMR signals

# Spatial encoding and field gradients

Gradient pulse causes Larmor frequency to become a function of position

If a delay is introduced between application of gradient pulse and readout, then the phase of the nuclear precession will become a function of position

Exploit these features to:

- Encode  $x$  position into  $k_x$  via “frequency encoding”
- Encode  $y$  position into  $k_y$  via “phase encoding”

Remember, gradient pulses  $G_i$  are such that:

$$B_z(x, y, z, t) = B_0 + xG_x(t) + yG_y(t) + zG_z(t)$$

$G_x = \frac{\partial B_z}{\partial x}$ ; i.e. a magnetic-field gradient in  $x$  direction

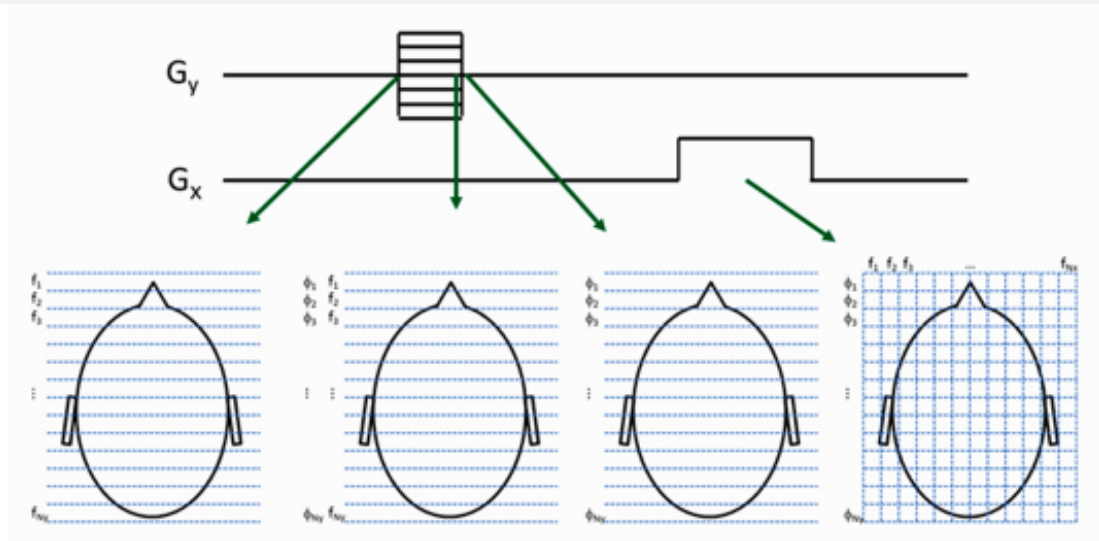
**magnetic field  $xG_x$  is in the  $\hat{k}$  direction**

$G_y = \frac{\partial B_z}{\partial y}$ ; i.e. a magnetic-field gradient in  $y$  direction

**magnetic field  $yG_y$  is in the  $\hat{k}$  direction**



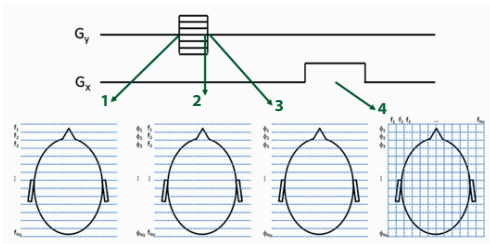
# Conversion of field gradient into $k$ space



# Conversion of field gradient into $k$ space

Example:

phase encode  $y$ ,  
frequency encode  $x$

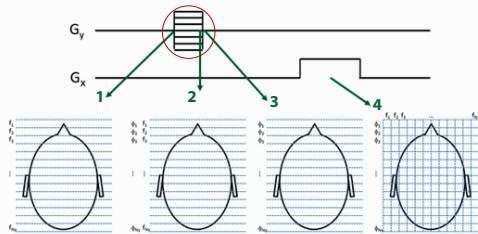


- ① At start of phase encoding-pulse, spins are in phase.  $G_y$  causes Larmor frequency to be function of  $y$ :  $\nu = f(y)$
- ② At end of phase encoding pulse, phase of precession,  $\phi$ , has become a function of  $y$ , i.e.  $\phi \rightarrow \phi(y)$
- ③ As time passes, phase dependence on  $y$  is preserved, i.e.  $\phi = \phi(y)$
- ④ Gradient pulse  $G_x$  causes Larmor frequency to become a function of  $x$ . Result is that  $y$ -position information is encoded in  $\phi = \phi(y)$  and  $x$ -position information is encoded in  $\nu = f(x)$

# Spatial encoding gradient pulses part of pulse sequence

Example:

phase encode  $y$ ,  
frequency encode  $x$



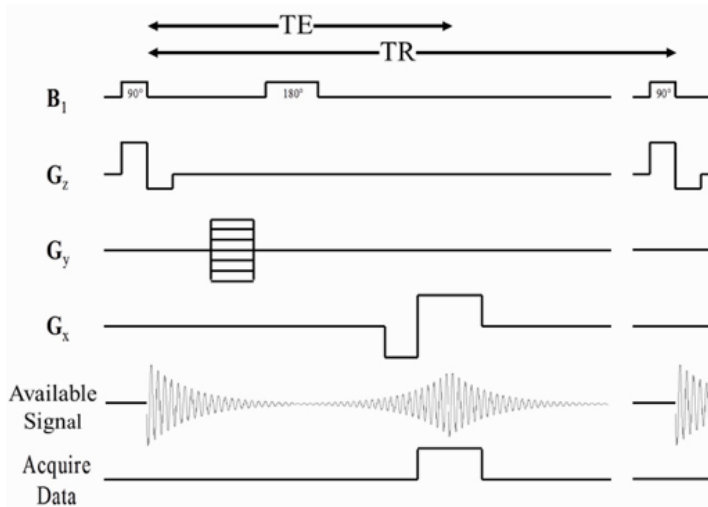
Phase- and frequency-encoding pulses part of a longer pulse sequence that repeats with period TR

At each repeat the amplitude of  $G_y$ , the phase-encoding pulse, has a different amplitude (as indicated on the figure)

For example:

- 1<sup>st</sup> iteration of sequence:  $G_y = 0$ ;
- 2<sup>nd</sup> iteration of sequence:  $G_y = +\eta$ ;
- 3<sup>rd</sup> iteration of sequence:  $G_y = -\eta$ ;
- ...

## Example pulse sequence



Example of spin-echo pulse sequence

Data is acquired at spin-echo time as shown

Combination of phase and frequency encoding pulses and repetition to obtain  $N_y$  data points completes ones transverse slice

## Section 2

# Lecture summary

# Summary

Slice-selective excitation achieved using magnetic field gradient such that Larmor frequency becomes function of position

Spatial information is encoded into net magnetisation in  $k$ -space

2D Fourier transform used to transform image in  $k$  space to image in coordinate space

$k$ -space image in slice is obtained using frequency and phase encoding

Pulse sequence is repeated to collect data for all  $N_x \times N_y$  pixels of image