

Evaluation of the proton fluxes captured and transmitted by plasma lenses (and solenoids) in LhARA

Titus-Stefan Dascalu

LhARA fortnightly Meeting
25th April 2023

Motivation

1. No previous investigation of the capability of the front-end LhARA beam-line to capture protons [without collimators](#)
 - No direct comparison of the [beam losses](#) for the case of using [plasma lenses vs. solenoids](#)
2. Significant impact of [vacuum nozzle](#) on [transmission](#) of protons from source to first lens
 - Previous optimisation of energy and momentum collimators, but not of the nozzle
3. [Full energy spectrum](#) of more realistic [protons](#) produced by [TNSA](#) available
 - from 3-D particle-in-cell simulations (WP2)

Studies presented here are based on the **baseline design of LhARA** Stage 1 (Sep 2022)

September 23, 2022

CCAP-TN-11 Issue 1

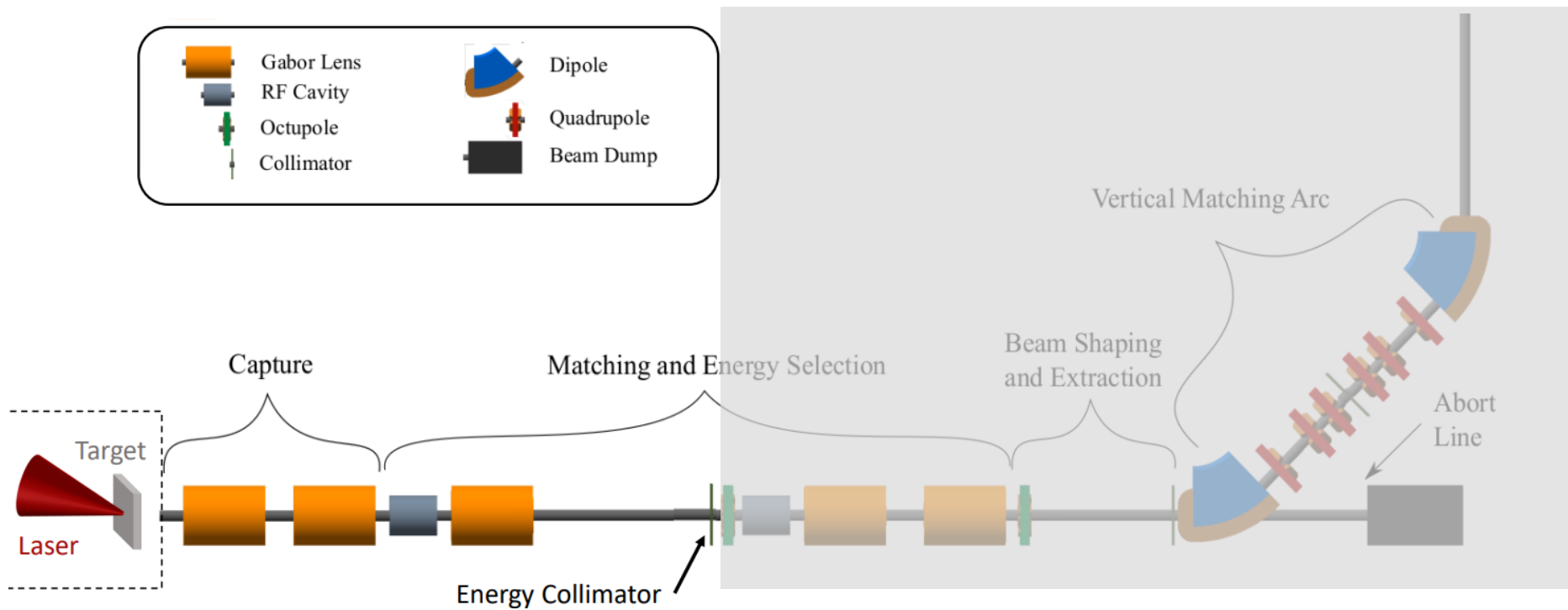


Baseline for the LhARA design update

The LhARA collaboration

- 1. Proton capture efficiency (without beam collimators)**
2. Protons loss due to beam collimators

Model only the front-end of LhARA Stage 1



- Protons tracked with GPT
 - first 5 cm from source without space-charge (due to comoving electrons)
 - with space-charge for the rest of the beam-line
- Plasma lenses modelled as field maps



Protons at the source

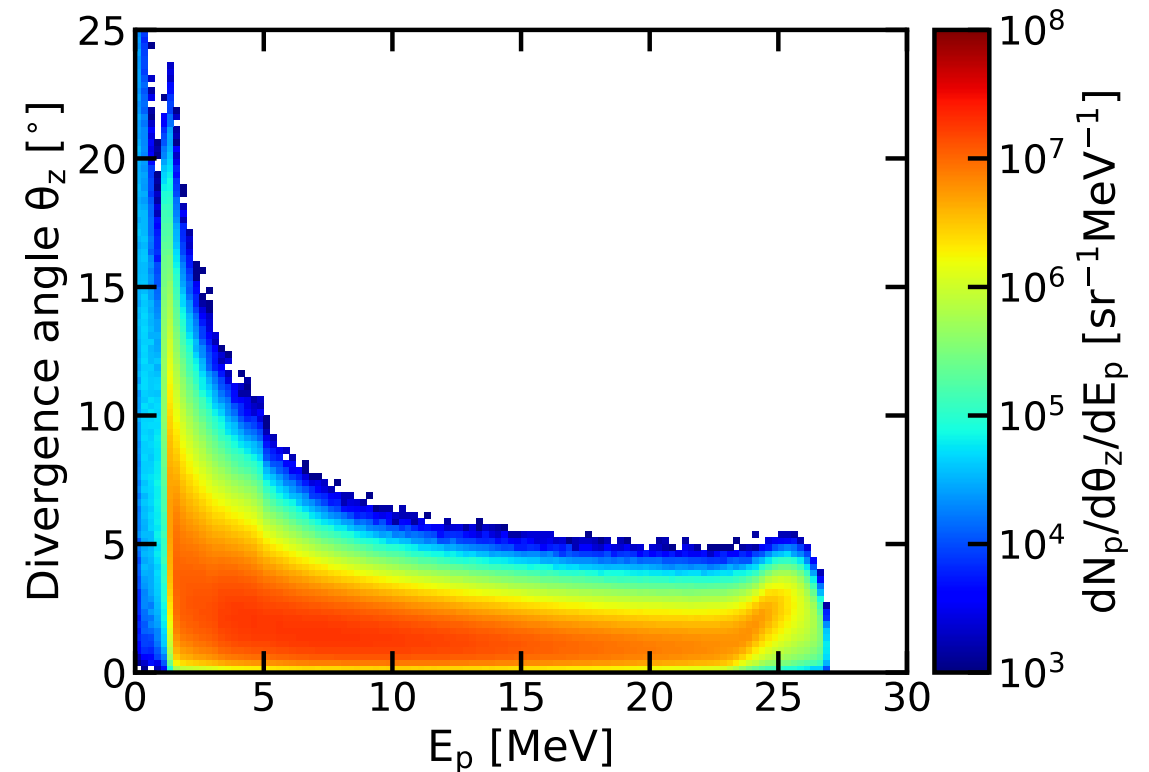
- 3-D particle-in-cell simulation of SCAPA-like laser hitting a solid tape target
 - TNSA regime
 - work of E. Boella described in the ITRF/LhARA 6M Progress Report

Key figures

full energy spectrum $\sim 10^{10}$ protons

nominal energy band $\sim 2 \times 10^9$ protons
(15 MeV \pm 2%)

RMS divergence angle $\sim 1^\circ$
(around nominal energy)

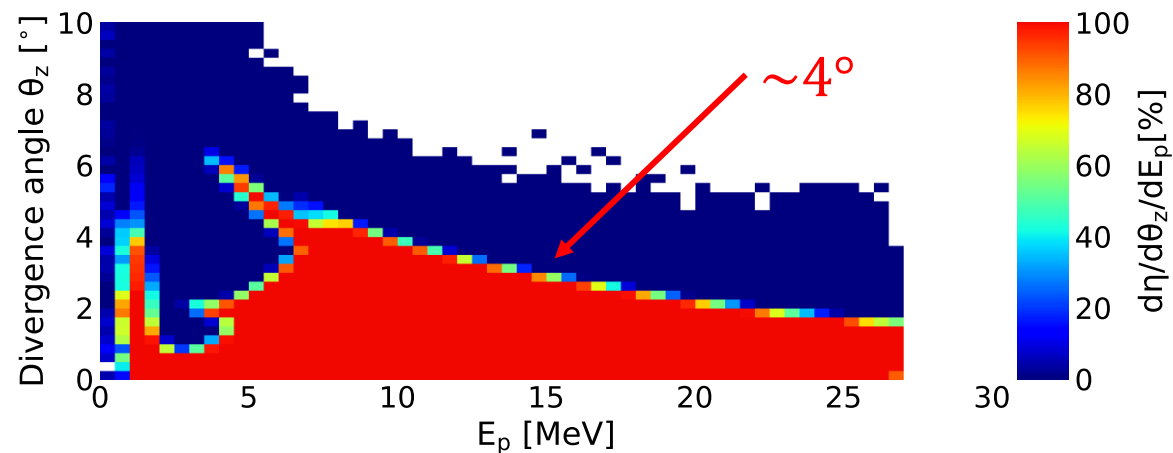
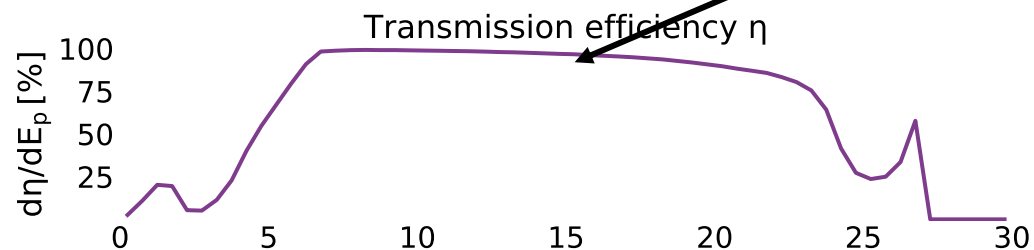


Transmission from source to first beam focus

- No vacuum nozzle, no energy collimator
 - Protons lost in the beam-pipe + inner walls of the solenoids/lenses

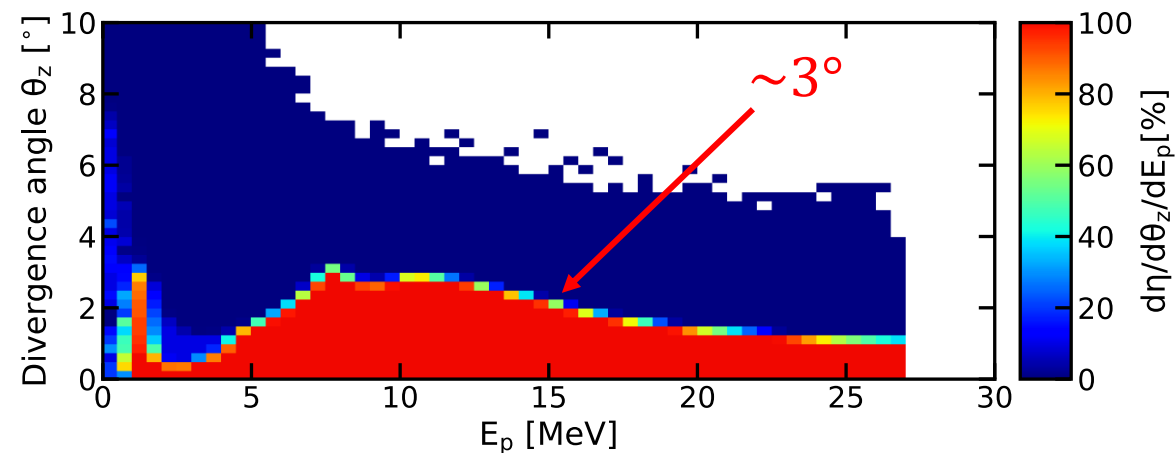
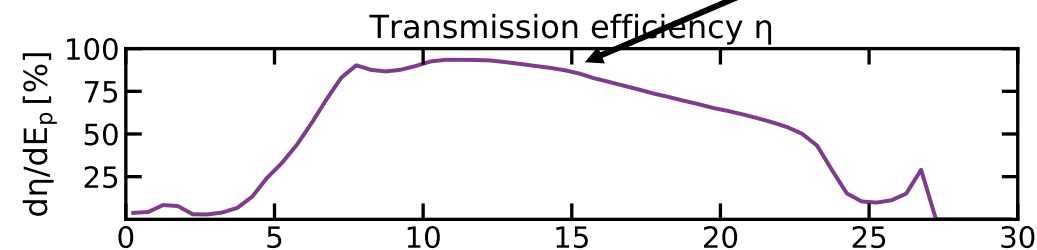
Plasma lenses

~98%



Solenoids

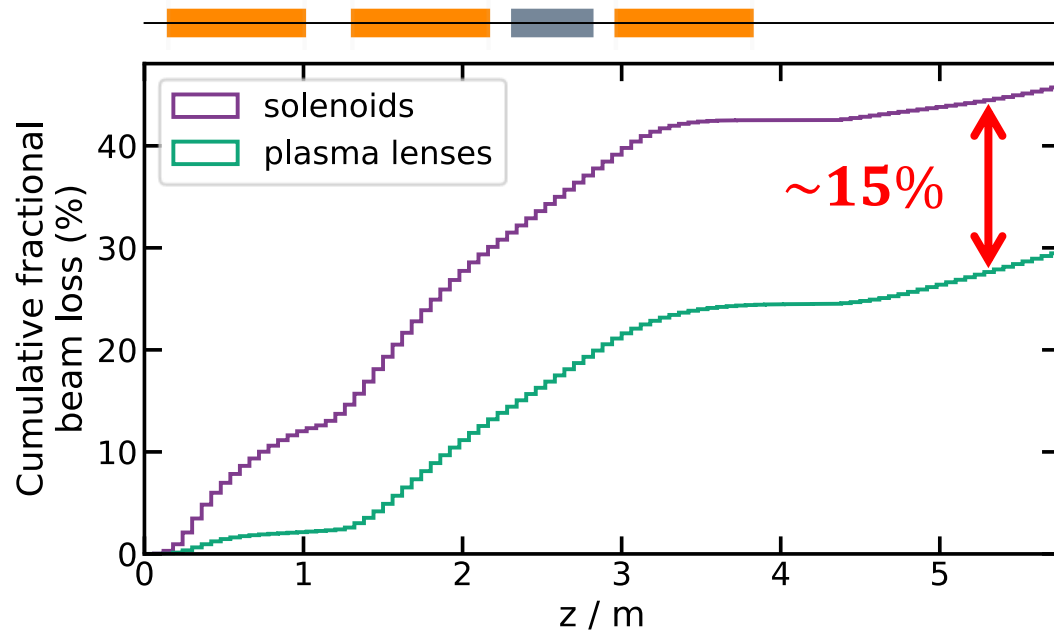
~85%



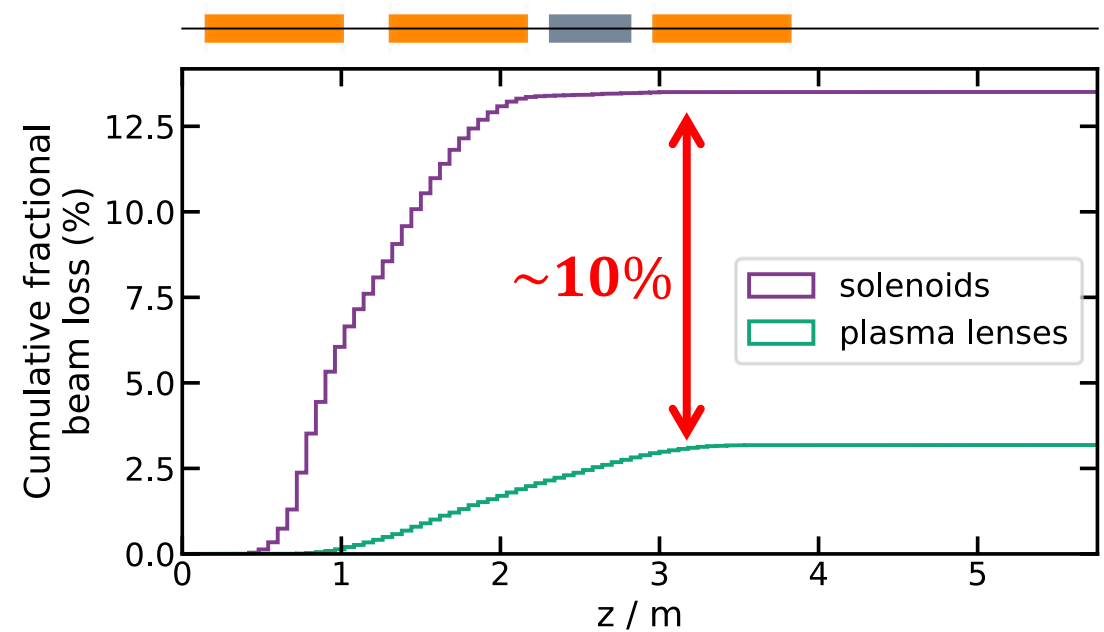
Beam loss integrated over angle and energy

- Beam loss calculated as fraction of number of protons produced at the source

- Full energy spectrum

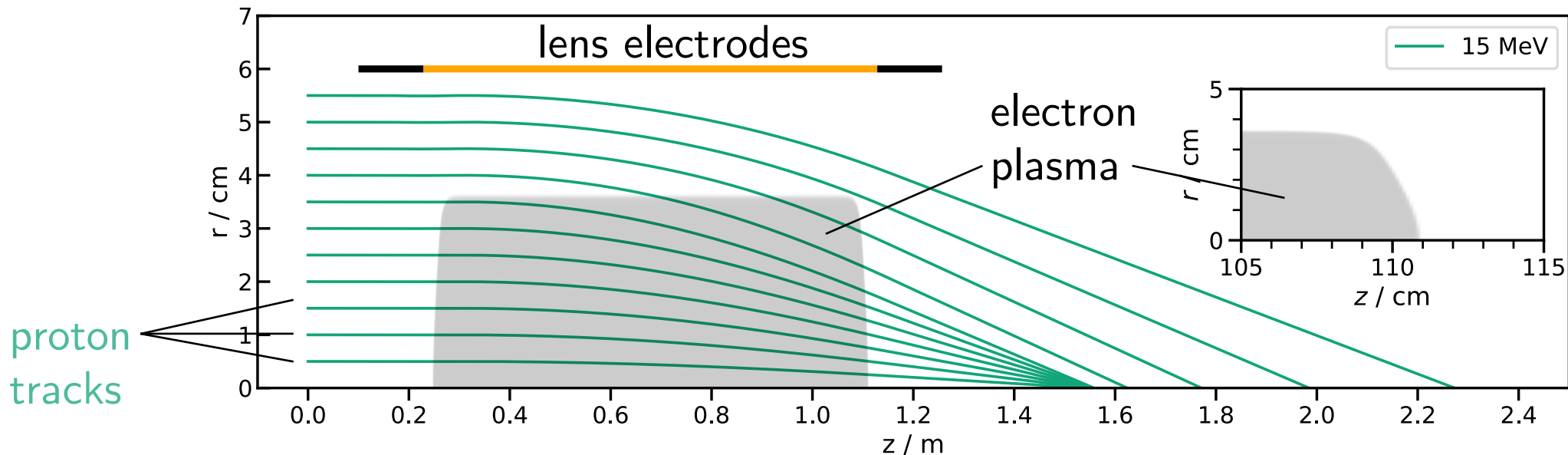


- Nominal energy band



Geometrical acceptance of Gabor lenses vs. solenoids

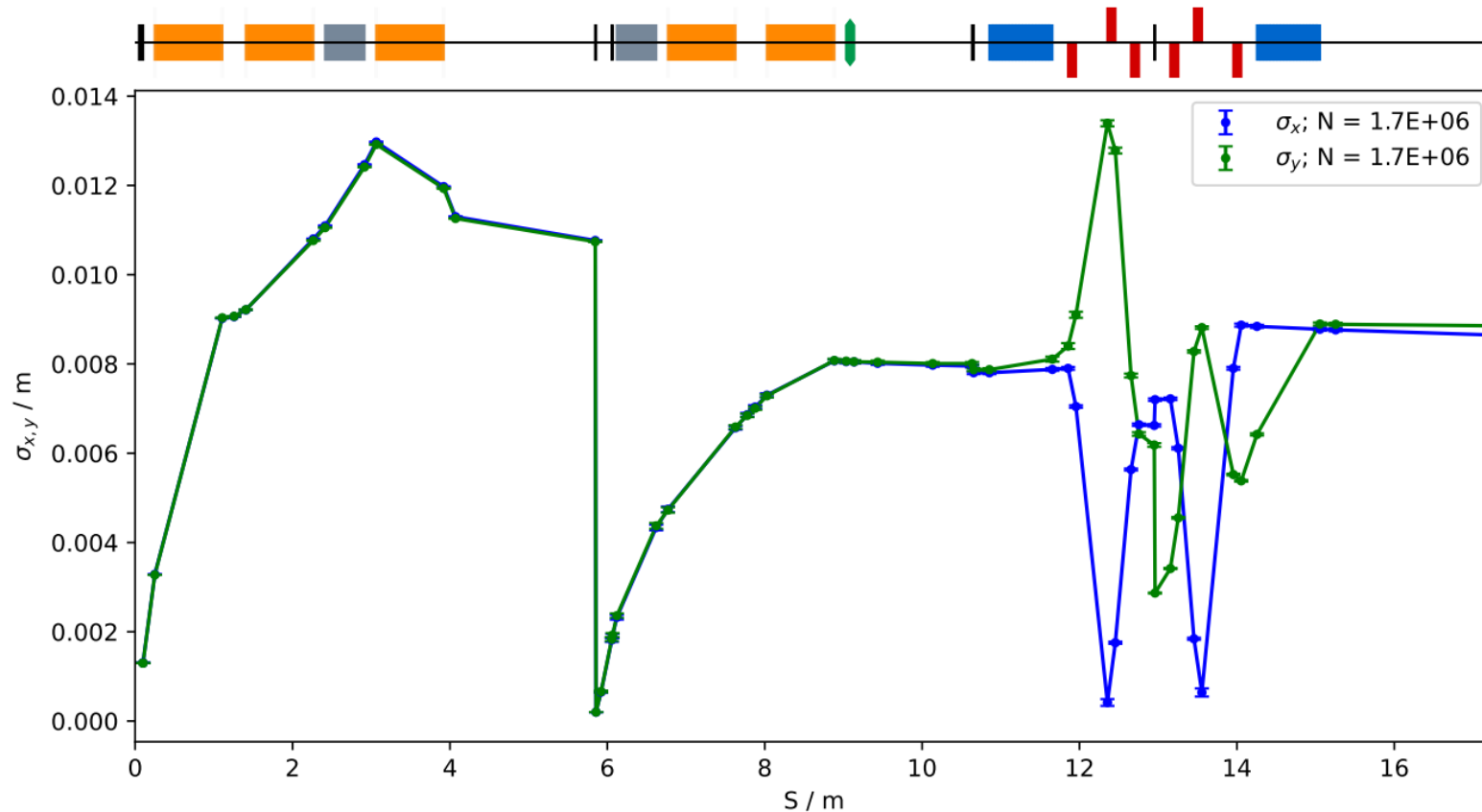
- Cost of large aperture of a solenoid is large overall size and mass
 - from preliminary design for normal-conducting solenoid for LhARA
- Plasma lenses generate additional focusing outside of the plasma
 - Non-linear focusing
- Plasma lenses should provide larger physical aperture
 - For diameter of the plasma identical to diameter of solenoid



1. Proton capture efficiency (without beam collimators)
- 2. Protons loss due to beam collimators**

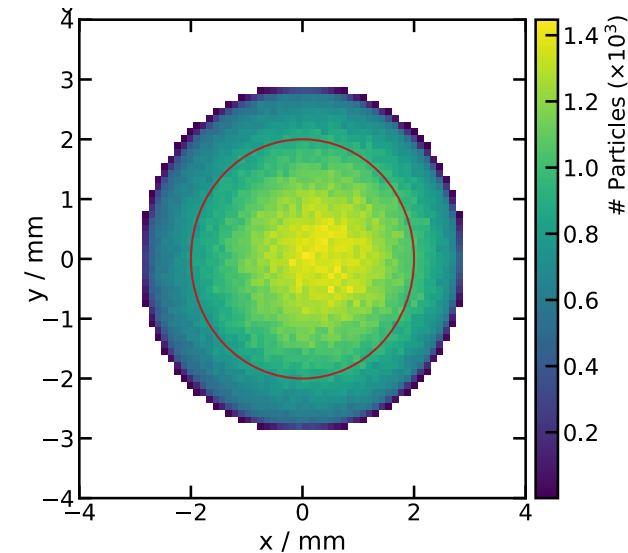
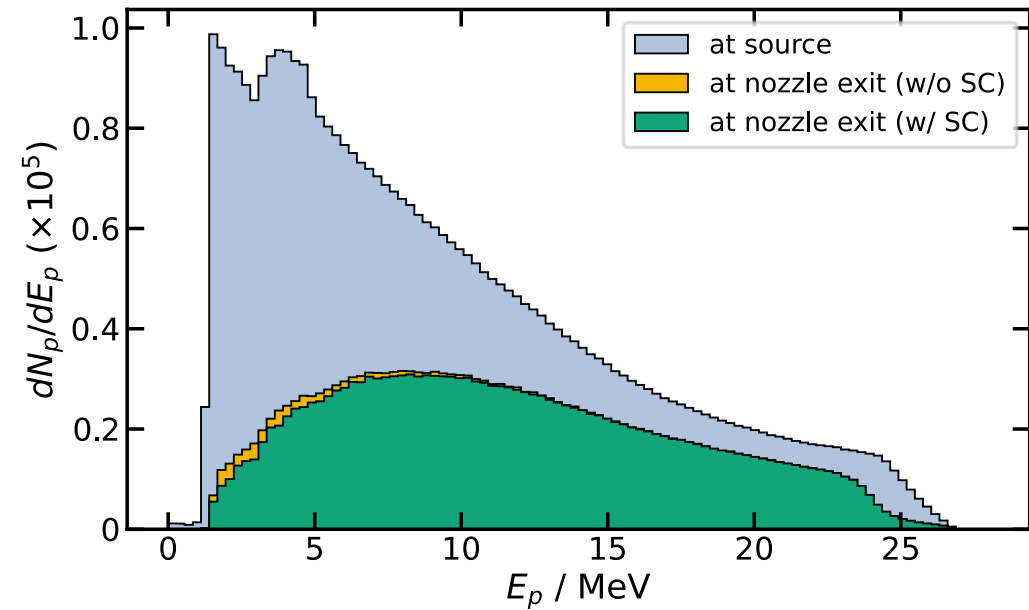
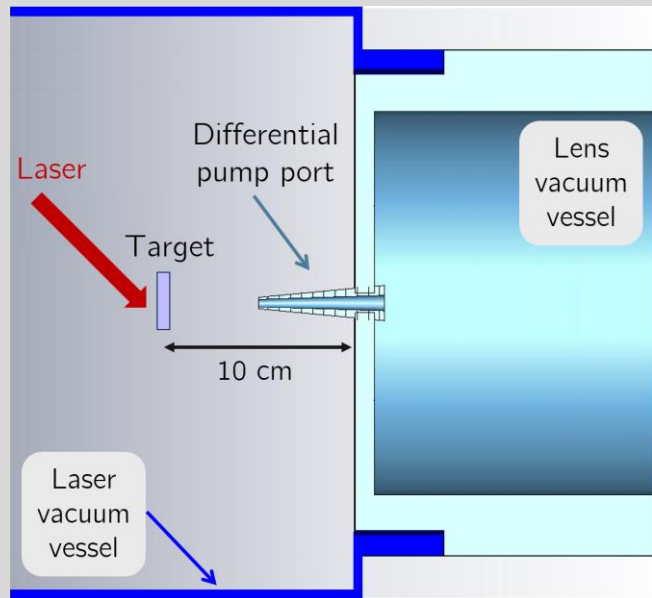
Model verification

- Protons with full energy spectrum
- Complete Stage 1 beam-line
- Hard-edge field maps for the Gabor lenses
- GPT for the first 10 cm, of which the last 5 cm with space-charge
- BDSIM for the rest of the beam-line, without space-charge



First significant reduction in proton flux

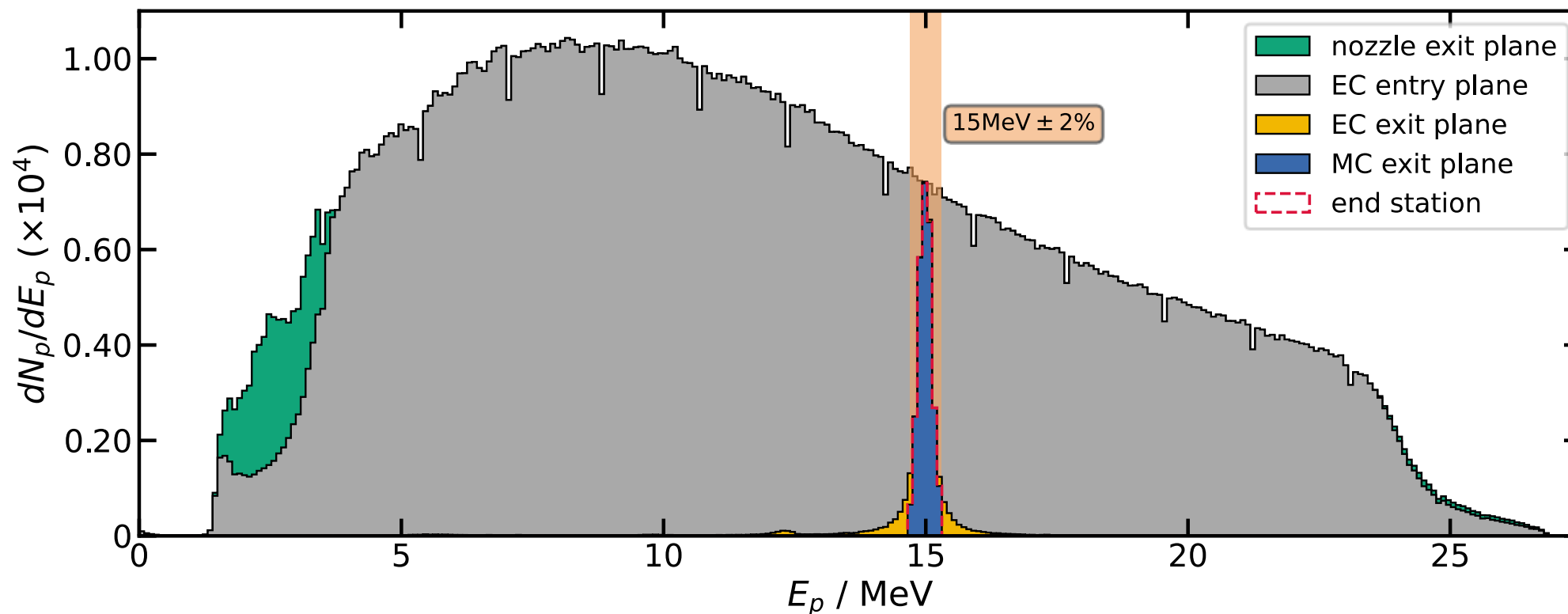
- At the interface between the source and the first lens
- Conical nozzle
 - 1° half-angle
 - 2 mm entrance radius



Validation of the energy-selection scheme

With solenoids

56% transmission from nozzle exit to end station



nozzle

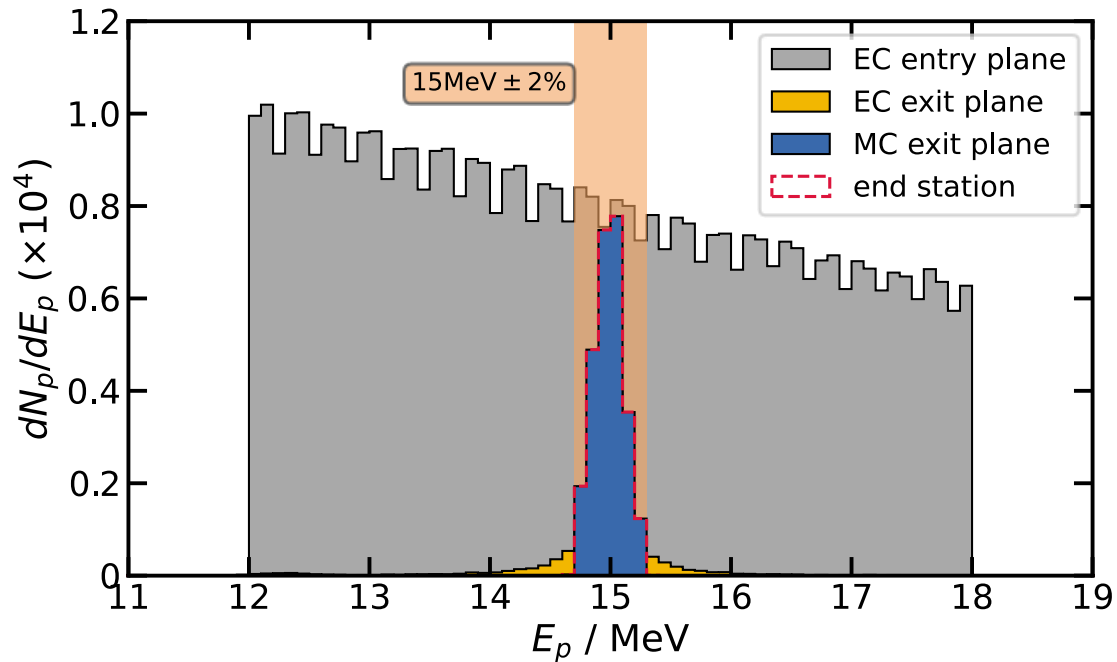
EC

MC

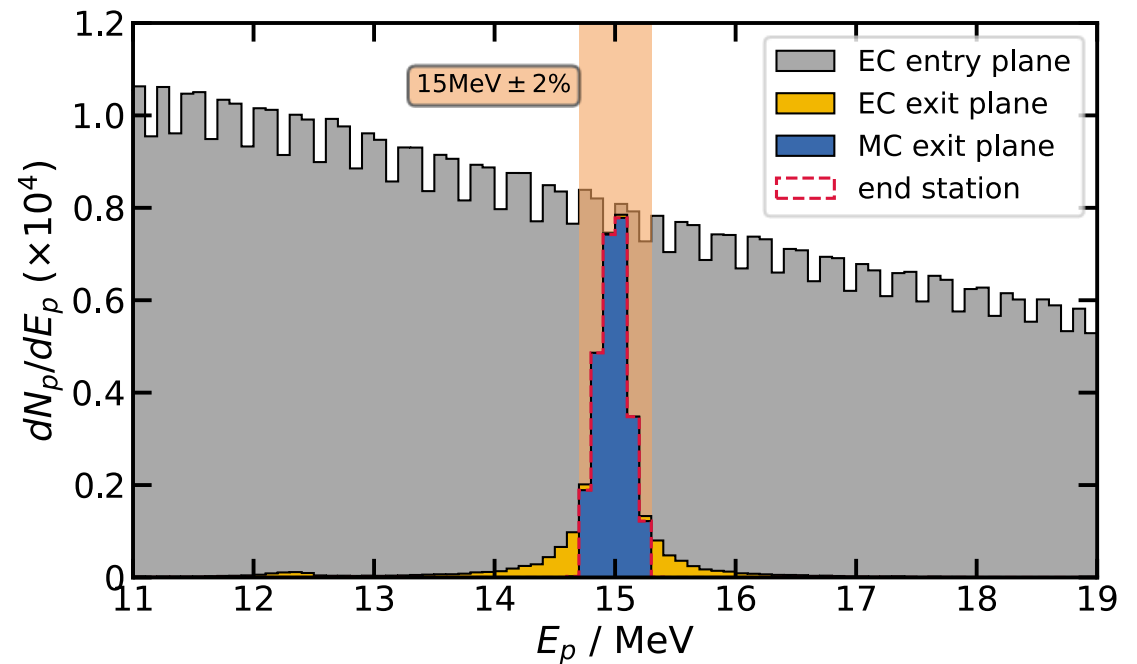


Identical energy selection with plasma/magnetic focusing

Plasma lenses



Solenoids



- 56% vs. 57% proton transmission within the nominal energy range from the nozzle

Conclusions

- In the absence of beam collimators, the capture section of LhARA transports a larger number of protons from the source when using plasma lenses compared to solenoids
- For the complete Stage 1 beam-line (baseline design, multiple collimators): identical fractions of protons reach the end station from the source with the use of solenoids or plasma lenses
 - within the nominal energy spread
- The superior capture efficiency of the plasma lenses compared to solenoids is suppressed by the limited angular acceptance of the nozzle situated downstream of the target
- Future optimisation of the nozzle should take into account
 - Beam-envelope size and divergence for protons within the nominal energy range from source
 - The full transport efficiency of the capture section of the LhARA beam-lines

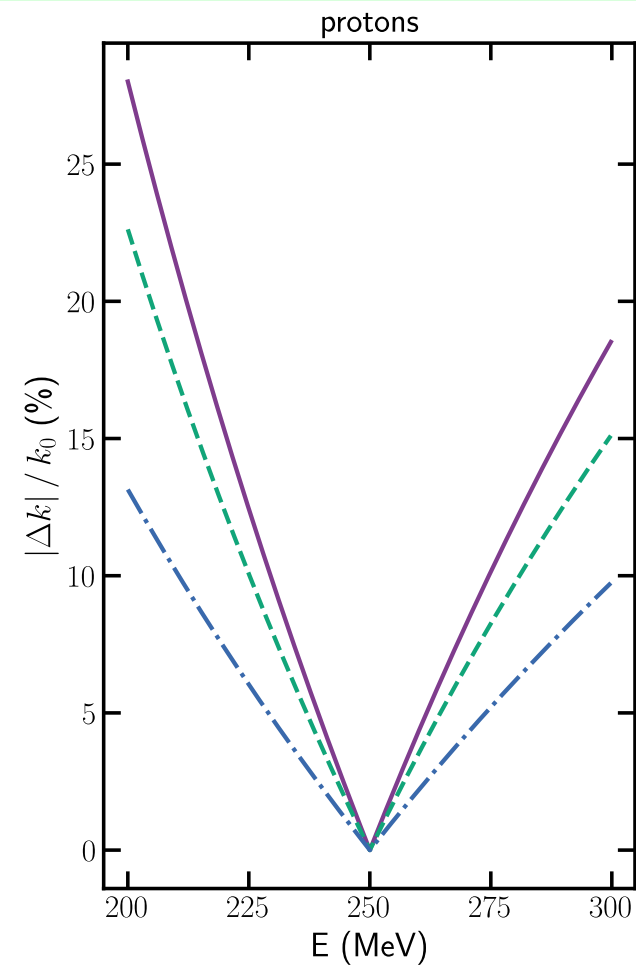
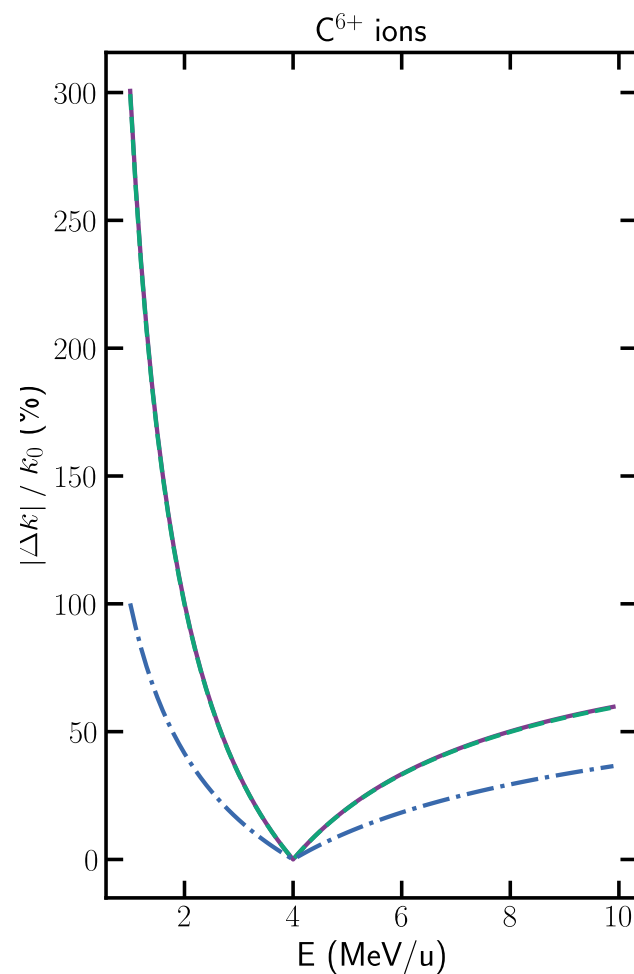
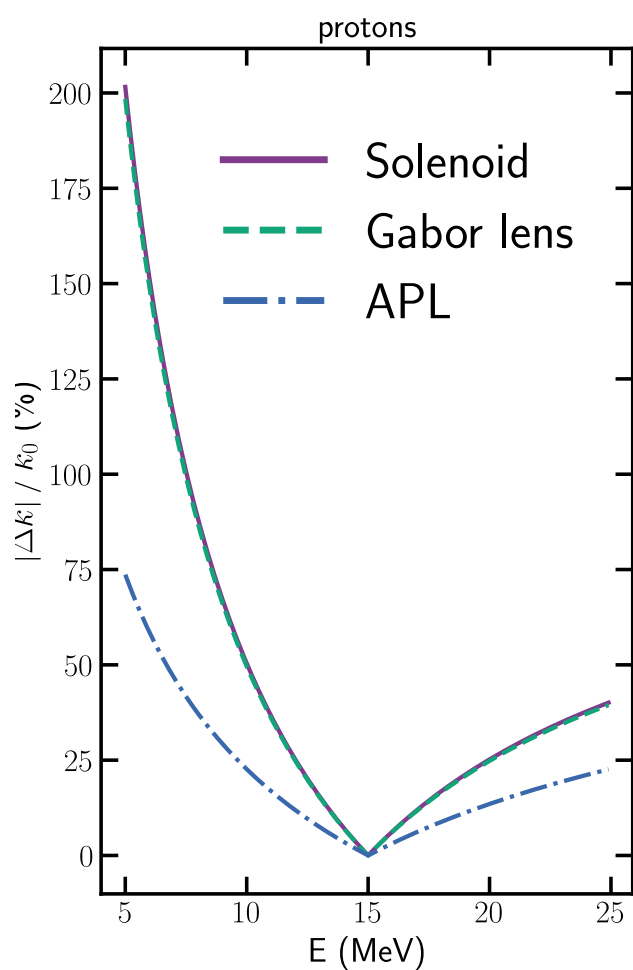
Back-up slides

Single-particle motion through a lens

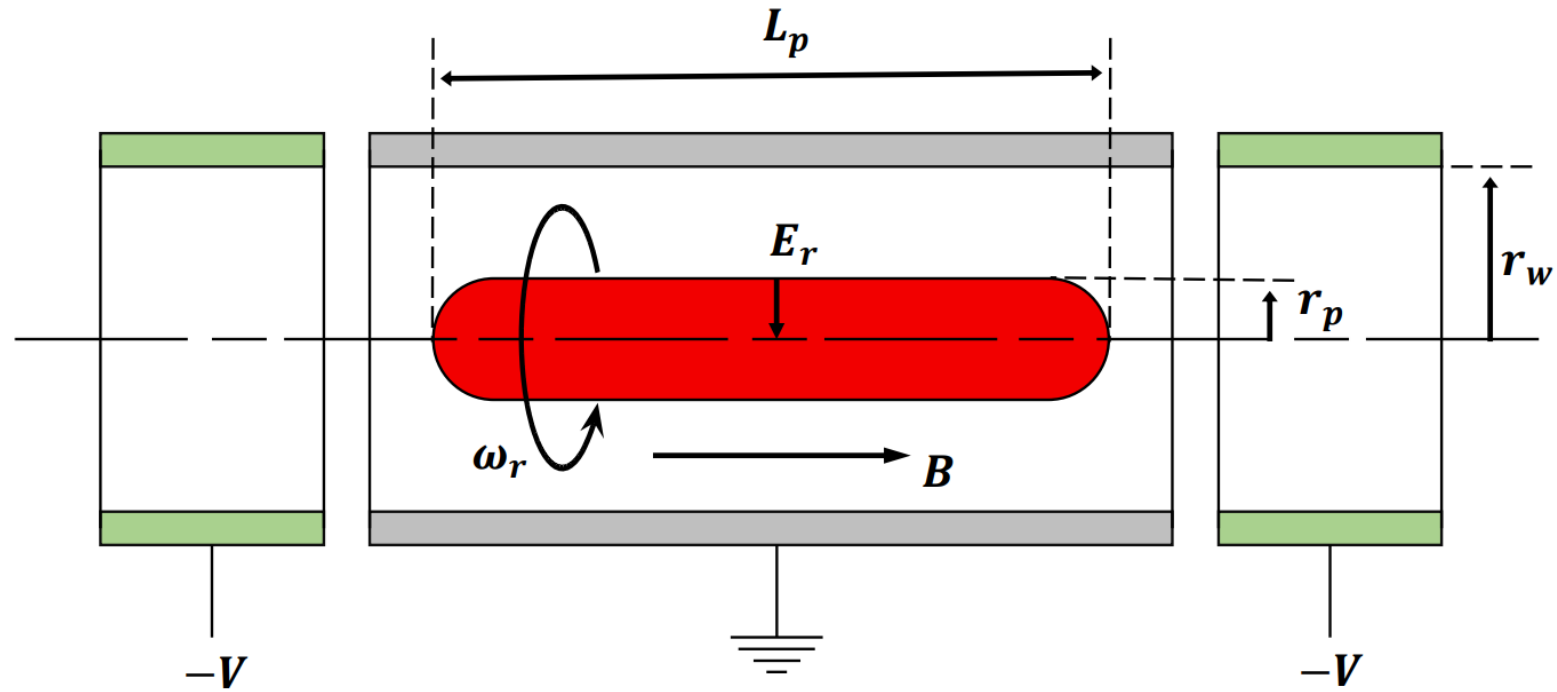
- Focusing strength parameter, $k \sim \frac{1}{f}$:

$$k_{GL} \sim \frac{\gamma_0}{P_0^2}$$

$$k_{\text{solenoid}} \sim \frac{1}{P_0^2}$$



Electron plasma (Gabor) lens

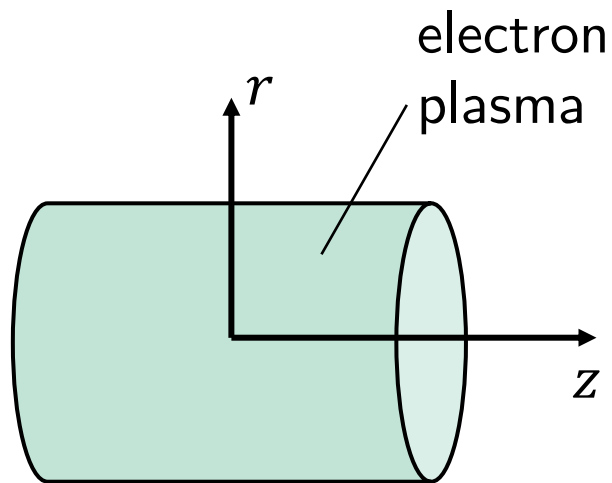


Field map of lenses with edge effects

[*] DOI: [10.1103/RevModPhys.87.247](https://doi.org/10.1103/RevModPhys.87.247)

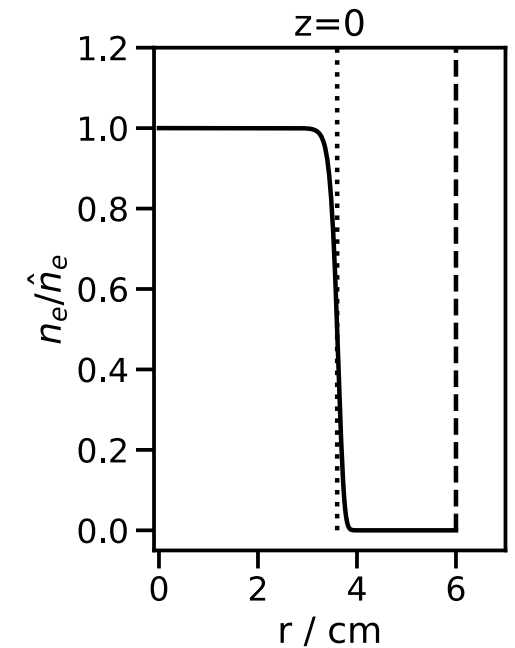
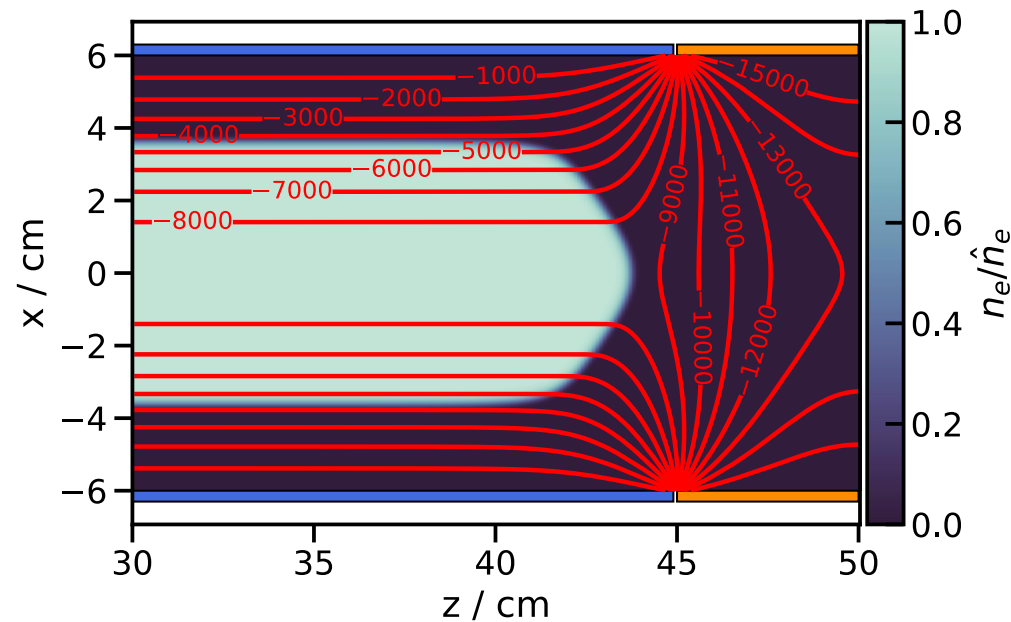
- Field map for each lens calculated separately
 - Plasma in global thermal equilibrium + rigid rotation
 - 2-D cylindrically symmetric numerical solution to Poisson-Boltzmann equations [1]

Hard-edge plasma (ideal lens)

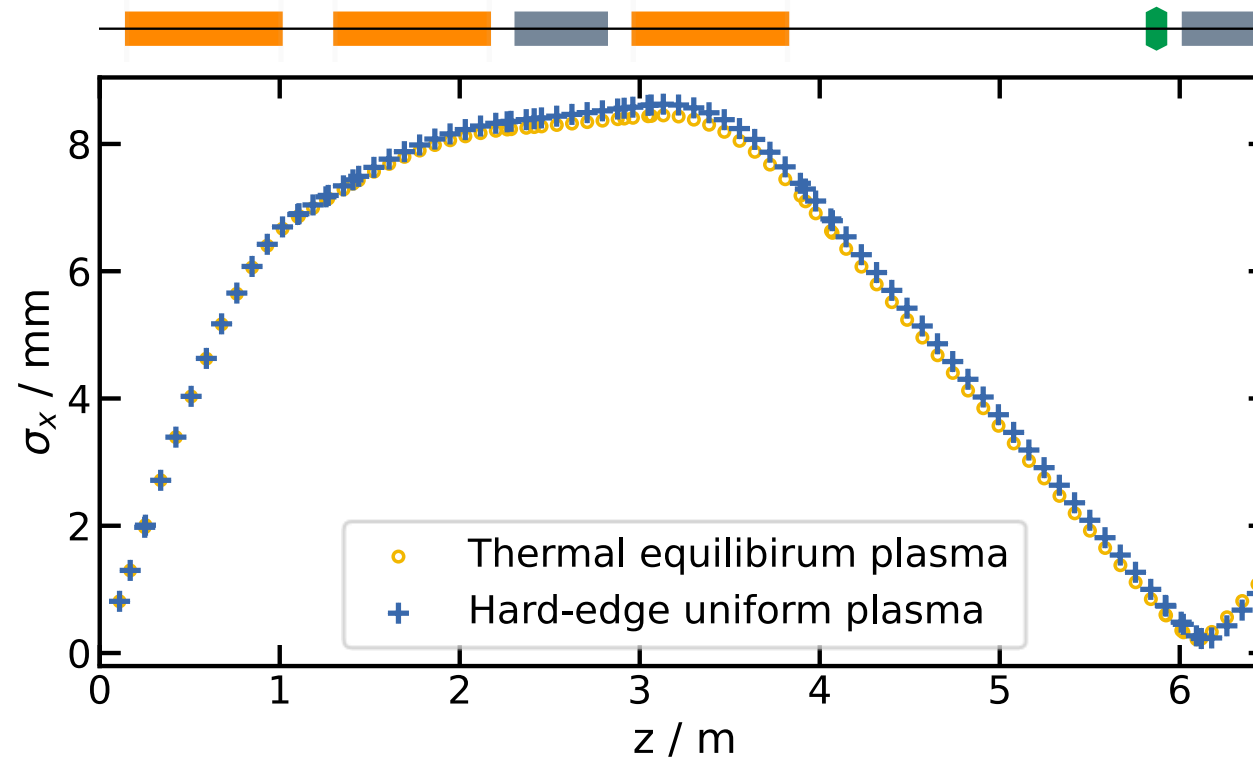


- uniform electron density
- infinitely long plasma

More realistic plasma shape



Contribution of edge-effects in beam-tracking



- Negligible differences between the two models of the lens
 - Preference for the hard-edge field map as it is much faster to generate (few minutes vs. several hours)