

Nuclear diagnostics and Magnetic Resonance Imaging

Week 1; Lecture 1; Section 3: Nuclear decay, revision

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Section 3

Revision of theory of nuclear decay

Radioactive decay

Observation:

“Activity”, A = number of disintegrations per second, decays exponentially with time

So, if there are N nuclei at time t :

$$N(t) = N(0) \exp(-\lambda t) \quad \dots \text{ and so } \dots \quad \frac{dN}{dt} = -\lambda N$$

λ is the decay constant. The “lifetime”, τ , and “half-life”, $t_{\frac{1}{2}}$, are related to λ by:

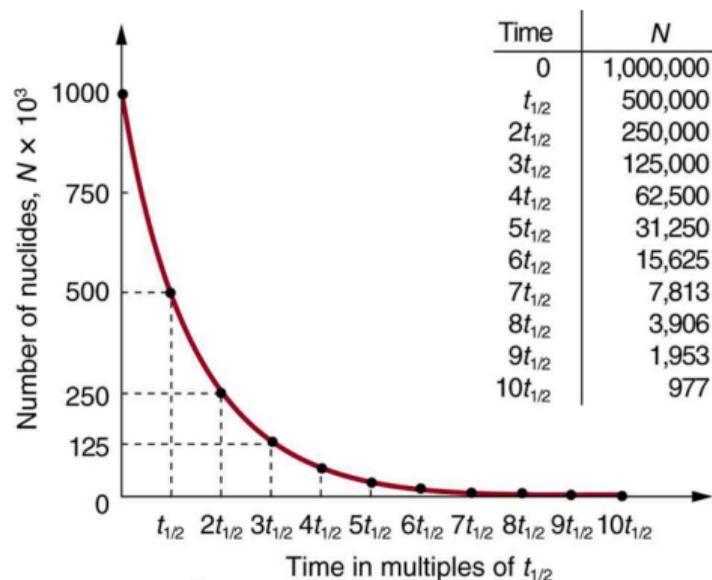
$$\tau = \frac{1}{\lambda} \quad \dots \text{ and } \dots \quad t_{\frac{1}{2}} = \frac{\ln(2)}{2} \tau = \frac{\ln(2)}{2} \frac{1}{\lambda}$$

The observation implies that the chance that a nucleus will decay per unit time is constant and is given by λ

Activity

Activity: $A = \left| \frac{dN}{dt} \right| = \gamma N$

- SI unity of activity: Becquerel (Bq):
 - 1 Bq = 1 decay per second
- Curie (Ci):
 - 1 Ci = 3.7×10^{10} Bq



<https://courses.lumenlearning.com/physics/chapter/31-5-half-life-and-activity/>

Branching ratio

Decay constant, λ , determines the decay rate. The 'lifetime', τ is defined to be:

$$\tau = \frac{1}{\lambda} \quad \dots \text{ and so } \dots \quad \lambda = \frac{1}{\tau}$$

The decay rate for the transition of X into Y may be calculated using "Fermi's Golden Rule":

$$\lambda_{X \rightarrow Y} = \eta |M_{X \rightarrow Y}|^2 \rho_f$$

Where η is a constant and ρ_f is the density of final states. $M_{X \rightarrow Y}$ is the quantum-mechanical 'matrix element' for the transition $X \rightarrow Y$. Some radionuclides may decay via more than one route. For such nuclei:

$$\lambda_T = \lambda_{X \rightarrow Y} + \lambda_{X \rightarrow Z} + \dots = \sum_i \lambda_i$$

λ_T is the 'total decay rate' (sometimes referred to as 'total width'). The λ_i are the 'partial' decay rates, or 'partial widths'.

Branching ratios; an additional constraint

'Branching ratio' (BR): fraction of all decays that result in a particular final state:

$$BR = \frac{\lambda_{X \rightarrow Y}}{\lambda_T}$$

Decay chain may include beneficial radiation, suitable for imaging, and harmful radiation.

Example: ^{131}I decay — $^{131}\text{I}(e^-, \gamma)^{131}\text{Xe}$

- Some γ s in useful range for imaging, but
- e^- and low-energy γ s simply deposit dose.

Has application in therapy, e.g., thyroid tumours.

Not widely used today for imaging.

So, consider ^{123}I , which decays via EC.

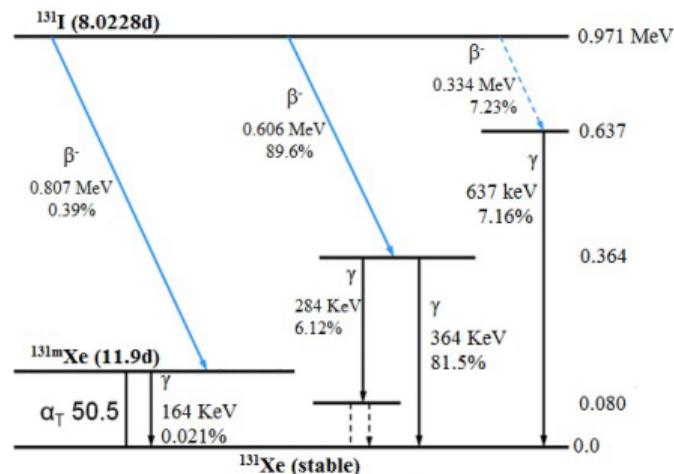
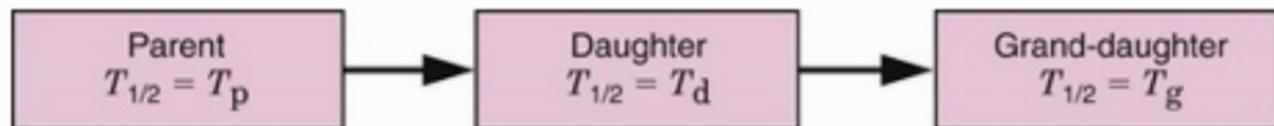


Figure from https://www.researchgate.net/publication/295919808_Radioiodine_I-131_for_Diagnosing_and_Treatment_of_Thyroid_Diseases

Parent-daughter decay chain



Branching ratio [Parent \rightarrow Daughter] = β

Rate of 'decay' of daughter nuclei:

$$\begin{aligned}\frac{dN_D}{dt} &= \lambda_P N_P \beta - \lambda_D N_D \\ &= \lambda_P N_{P0} \beta \exp(-\lambda_P t) - \lambda_D N_D\end{aligned}$$

i.e.:

$$\frac{dN_D}{dt} + \lambda_D N_D - \lambda_P \beta N_{P0} \exp(-\lambda_P t) = 0.$$

Solution:

$$N_D = \frac{\lambda_P}{\lambda_D - \lambda_P} \beta N_{P0} [\exp(-\lambda_P t) - \exp(-\lambda_D t)] + N_{D0} \exp(-\lambda_D t)$$

Or in terms of activation:

$$A_D = \frac{\lambda_D}{\lambda_D - \lambda_P} \beta A_{P0} [\exp(-\lambda_P t) - \exp(-\lambda_D t)] + A_{D0} \exp(-\lambda_D t) \quad (1)$$

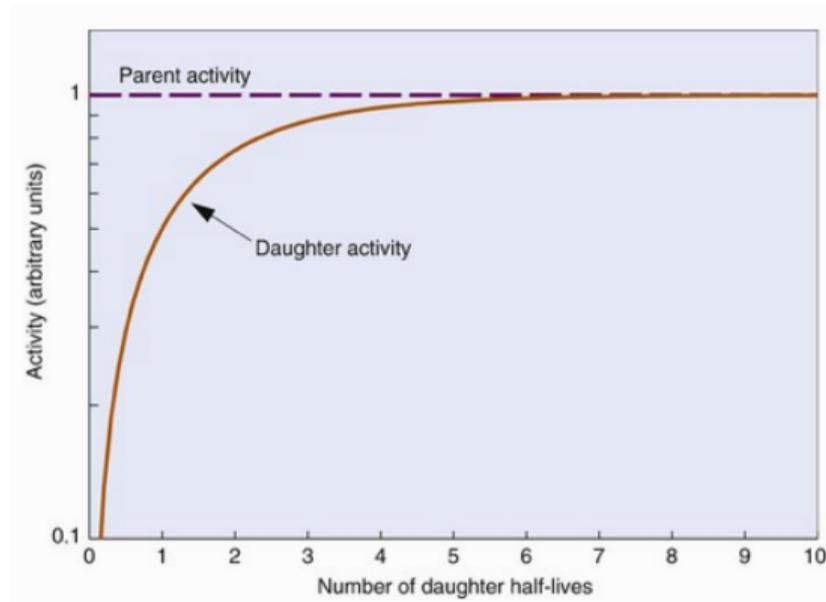
$T_P \gg T_D$; secular equilibrium

$$T_P \gg T_D \Rightarrow \frac{\lambda_P}{\lambda_D} \ll 1 \text{ and } \exp(-\lambda_P t) \sim 1.$$

So, equation 1 becomes:

$$A_D = \beta A_{P0} [1 - \exp(-\lambda_D t)] + A_{D0} \exp(-\lambda_D t)$$

If $A_{D0} = 0$ and $\beta = 1$, then the build up of N_D reaches 'secular equilibrium after 5–6 T_D .



$T_P > T_D$; transient equilibrium

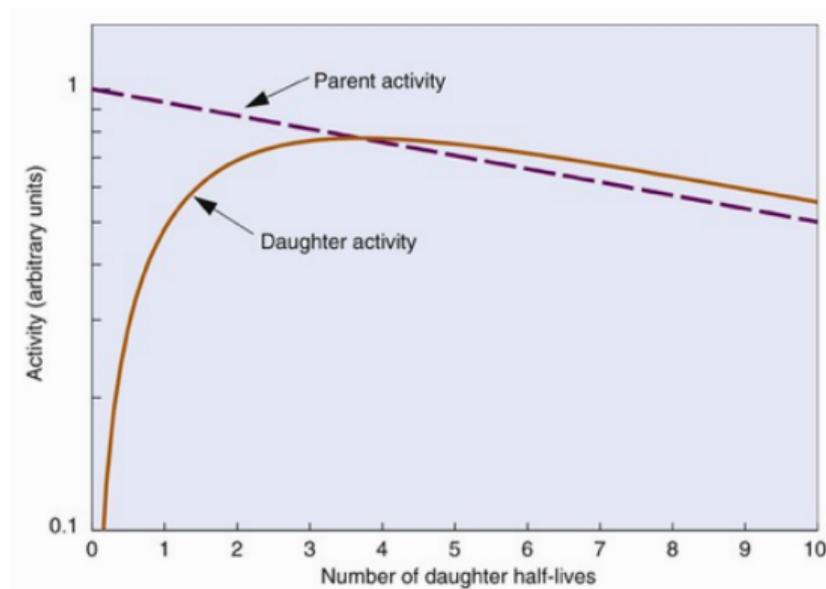
Transient equilibrium occurs at t_{eq} given by:

$$t_{\text{eq}} = \frac{\ln \left[\frac{\lambda_P}{\lambda_D} \right]}{\lambda_P - \lambda_D}$$

At this time the activity of the daughter is a maximum, so one may write:

$$t_{\text{max}} = t_{\text{eq}} = \frac{1.44 T_P T_D}{T_P - T_D} \ln \left[\frac{T_P}{T_D} \right]$$

If $A_{D0} = 0$, $\beta = 1$, and $T_D = 0.1 T_P$, then build up and decay of N_D reaches 'transient equilibrium' after $\sim 2.6 T_D$.



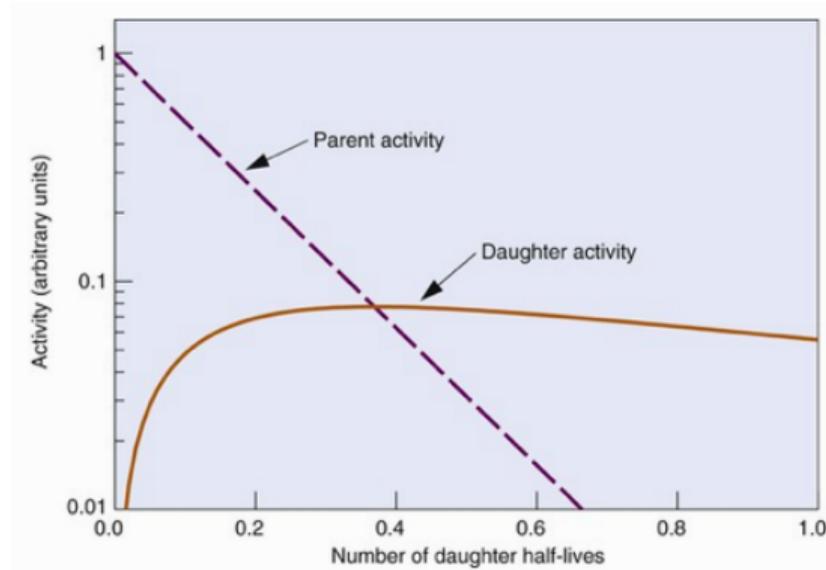
$T_P < T_D$; no equilibrium

Maximum activity of daughter is still given by:

$$t_{\max} = t_{\text{eq}} = \frac{1.44 T_P T_D}{T_P - T_D} \ln \left[\frac{T_P}{T_D} \right]$$

The daughter activity grows until t_{\max} and then decreases. The parent activity 'falls away' and therefore fails to replenish the daughter.

If $A_{D0} = 0$, $\beta = 1$, and $T_D = 10 T_P$, then build up of N_D reaches maximum activity at $\sim 0.26 T_D$.



Summary of section 3

Radioactive decay law:

$$\frac{dN}{dt} = -\lambda N \exp(-\lambda t)$$

implies that chance that a nucleus will decay per unit time is λ

Branching ratio for decay $X \rightarrow Y$: $\text{BR} = \frac{\lambda_{X \rightarrow Y}}{\lambda_T}$

Can solve for evolution of samples of parent, daughter, grand-daughter, etc. in decay chain and distinguish between:

- Secular equilibrium for $T_P \gg T_D$
- Transient equilibrium for $T_P > T_D$
- No equilibrium for $T_P < T_D$