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Magnetic Resonance Imaging Week 4; Lecture 8; Section 2: Rotating the magnetisation

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Section 2

Rotating the magnetisation

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First, a static example



Consider magnetisation **M** parallel to z axis and **B** parallel to the x axis, as shown

Torque, $\mathbf{M} \times \mathbf{B}$, is therefore parallel to the y axis

Net result is that **M** will precess around the x axis towards the y axis

This is what is done in MRI

Rotating the magnetisation vector in MRI; principle

Main field, \mathbf{B}_0 , produced with solenoid



Induces magnetisation \boldsymbol{M} parallel to \boldsymbol{B}_0

To rotate **M** away from **B**₀ require magnetic field in transverse (x, y) plane

Call the field in the x, y plane **B**₁; can be produced with a variety of coil arrangements, e.g. dipole or, more efficient, a "bird cage"



To cause **M** to precess require that **M** oscillates at the Larmor frequency, ω . I.e. require RF magnetic field **B**₁

Rotating the magnetisation vector in MRI; mathematics

Take **B**₁ to be "plane polarised" in x, y such that $B_{1_x} = B_1 \cos(\omega t + \alpha)$ and $B_{1_y} = B_1 \sin(\omega t + \beta)$; α and β are phases

 \mathbf{B}_1 can be rewritten in terms of two circularly polarised fields:



Rotating the magnetisation vector in MRI

One of the two counter rotating fields will rotate in the same direction as the nuclear precession

In the frame that is co-rotating with the precession of the net magnetisation vector the magnetic field will appear stationary in the transverse (x, y) plane. Call the co-rotating field B_1^+

 B_1^+ is equal to either $B_{1_{ac}}$ or B_{1_c} depending on the direction of \mathbf{B}_0

The stationary field will therefore cause **M** to precess about a rotating axis in the (x, y) plane

The net result is that **M** can be rotated into the x, y plane where it will continue to precess

The precession of \mathbf{M} in the x, y plane gives a detectable RF signal

Rotating the magnetisation vector in MRI



M is initially parallel to \mathbf{B}_0

(a) Laboratory Frame of Reference

(b) **Rotating Frame of Reference**

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The flip angle

The flip angle, α , is proportional to the magnitude and duration of the RF pulse:

 $\alpha = \gamma B_1 t_P$

where t_P is the duration of the RF pulse

 90° pulse rotates magnetisation into transverse plane where it continues to precess

Effect of 90° RF Pulse



Example: calculating the duration of a 90° pulse

RF transverse magnetic field pulse is applied to rotate ${\bf M}$

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The magnitude of B_1 is 10 \muT (i.e. 10<sup>-5</sup> T)
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At what rate with the $\boldsymbol{\mathsf{M}}$ rotate away from the $\boldsymbol{\mathsf{B}}_0$ axis?

How long will it take for the flip angle to reach 90° ?

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Example: calculating the duration of a 90° pulse

Half an answer . . .

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The magnitude of B_1 is 10 \muT (i.e. 10<sup>-5</sup> T)
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At what rate with the **M** rotate away from the **B**₀ axis? It will rotate at the Larmor frequency, f_1 arising from the field B_1 , i.e. $f_1 = \gamma B_1$

How long will it take for the flip angle to reach 90°? The angle can be obtained by solving the equation:

$$rac{1}{4}=\gamma B_1 t_P^{90^\circ} ext{ for } t_P^{90^\circ}$$
 or $rac{\pi}{2}=\gamma B_1 t_P^{90^\circ}$

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Summary of section 2

Net magnetisation of ¹H spins caused to rotate using plane-polarised, time-varying magnetic field in the x, y plane

Precession of rotated net-magnetisation vector gives rise to RF signal which can be detected

Measurement of the RF signal from the precession of rotated net-magnetisation vector is the basis of MRI

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