

Physics of Medical Imaging and Radiotherapy

Magnetic Resonance Imaging

Lecture 1; Section 3: Classical derivation of the Larmor equation

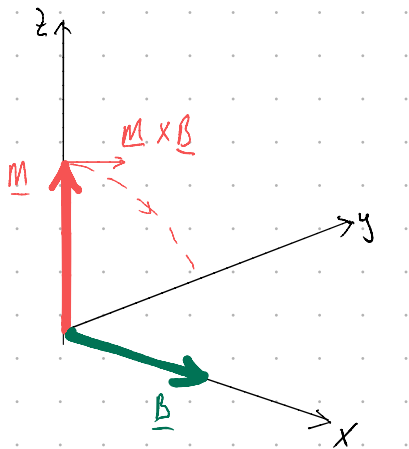
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Section 1

Rotating the magnetisation

First, a static example



Consider magnetisation \mathbf{M} parallel to z axis and \mathbf{B} parallel to the x axis, as shown

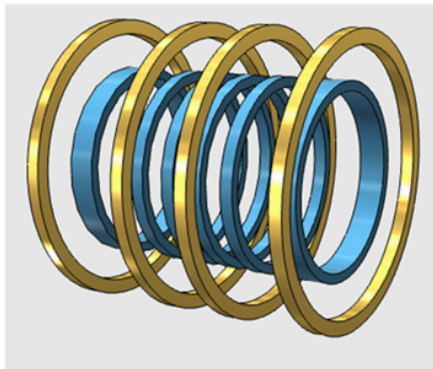
Torque, $\mathbf{M} \times \mathbf{B}$, is therefore parallel to the y axis

Net result is that \mathbf{M} will precess around the x axis towards the y axis

This is what is done in MRI ...

Rotating the magnetisation vector in MRI; principle

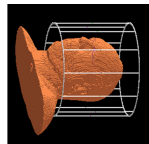
Main field, \mathbf{B}_0 , produced with solenoid



Induces magnetisation \mathbf{M} parallel to \mathbf{B}_0

To rotate \mathbf{M} away from \mathbf{B}_0 require magnetic field in transverse (x, y) plane

Call the field in the x, y plane \mathbf{B}_1 ; can be produced with a variety of coil arrangements, e.g. dipole or, more efficient, a “bird cage”



To cause \mathbf{M} to precess require that magnetic field oscillates at the Larmor frequency, ω .

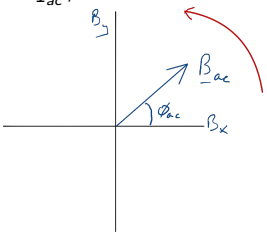
... i.e. require RF magnetic field \mathbf{B}_1

Rotating the magnetisation vector in MRI; mathematics

Take \mathbf{B}_1 to be “plane polarised” in x, y such that $B_{1x} = B_1 \cos(\omega t + \alpha)$ and $B_{1y} = B_1 \sin(\omega t + \beta)$; α and β are phases

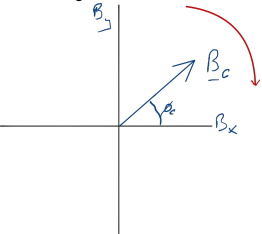
\mathbf{B}_1 can be rewritten in terms of two circularly polarised fields:

$\mathbf{B}_{1_{ac}}$; anti-clockwise



$B_{1_{ac}} = \frac{B_1}{2}; \phi_{ac} = \omega t + \alpha'$

\mathbf{B}_{1_c} ; clockwise



$B_{1_c} = \frac{B_1}{2}; \phi_c = \omega t + \beta'$

Rotating the magnetisation vector in MRI

One of the two counter rotating fields will rotate in the same direction as the nuclear precession

Either $B_{1_{ac}}$ or B_{1_c} will appear stationary in the plane transverse to \mathbf{B}_0 in the frame that is co-rotating with the precession of the net magnetisation vector. Call the co-rotating field B_1^+

B_1^+ is equal to either $B_{1_{ac}}$ or B_{1_c} depending on the direction of \mathbf{B}_0

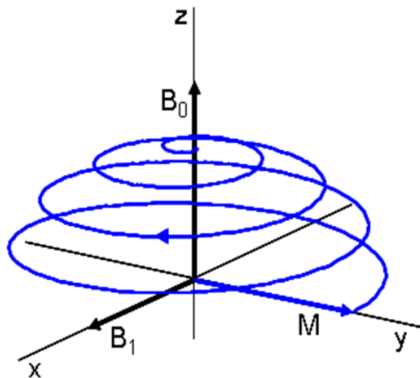
The field stationary in the rotating frame will therefore cause \mathbf{M} to precess about a rotating axis in the transverse plane

The net result is that \mathbf{M} can be rotated into the x, y plane where it will continue to precess

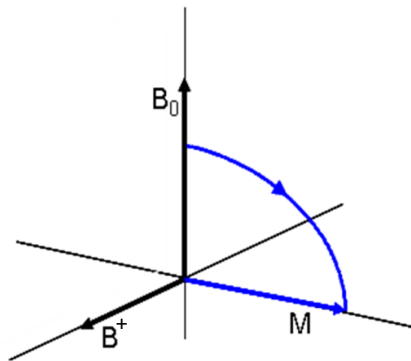
The precession of \mathbf{M} in the x, y plane gives a detectable RF signal

Rotating the magnetisation vector in MRI

M is initially parallel to B_0



(a) *Laboratory Frame of Reference*



(b) *Rotating Frame of Reference*

The flip angle

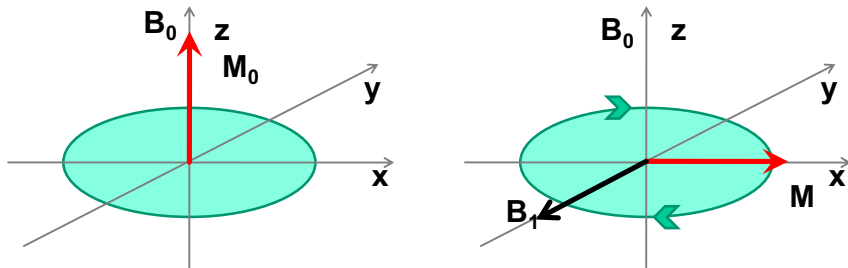
The flip angle, α , is proportional to the magnitude and duration of the RF pulse:

$$\alpha = \gamma B_1 t_P$$

where t_P is the duration of the RF pulse

90° pulse rotates magnetisation into transverse plane where it continues to precess

Effect of 90° RF Pulse



Example: calculating the duration of a 90° pulse

RF transverse magnetic field pulse is applied to rotate \mathbf{M}

The magnitude of B_1 is $10 \mu\text{T}$ (i.e. 10^{-5} T)

At what rate will \mathbf{M} rotate away from the \mathbf{B}_0 axis?

How long will it take for the flip angle to reach 90° ?

Example: calculating the duration of a 90° pulse

Half an answer . . .

The magnitude of B_1 is $10 \mu\text{T}$ (i.e. 10^{-5}T)

At what rate will \mathbf{M} rotate away from the \mathbf{B}_0 axis?

It will rotate at the Larmor frequency, f_1 , arising from the field B_1 , i.e. $f_1 = \gamma B_1$

How long will it take for the flip angle to reach 90° ?

The angle can be obtained by solving the equation:

$$\frac{1}{4} = \gamma B_1 t_P^{90^\circ} \text{ for } t_P^{90^\circ}$$
$$\text{or } \frac{\pi}{2} = \gamma B_1 t_P^{90^\circ}$$

Summary of section 3

Net magnetisation of ^1H spins caused to rotate using plane-polarised, time-varying magnetic field in the x, y plane

Precession of rotated net-magnetisation vector gives rise to RF signal which can be detected

Measurement of the RF signal from the precession of rotated net-magnetisation vector is the basis of MRI